

# Introduction to Computational Electromagnetics

## Lecture 3: 2D Electromagnetic Scattering

### Part 2: The Magnetic Field Integral Equation

ELC 657 – Spring 2014

Department of Electronics and Communications  
Engineering  
Faculty of Engineering – Cairo University

# Outline

## 1 TM Polarization

- Problem Formulation
- MoM Procedure
- Sample Results

## 2 TE Polarization

- Problem Formulation
- MoM Procedure
- Sample Results

# Outline

## 1 TM Polarization

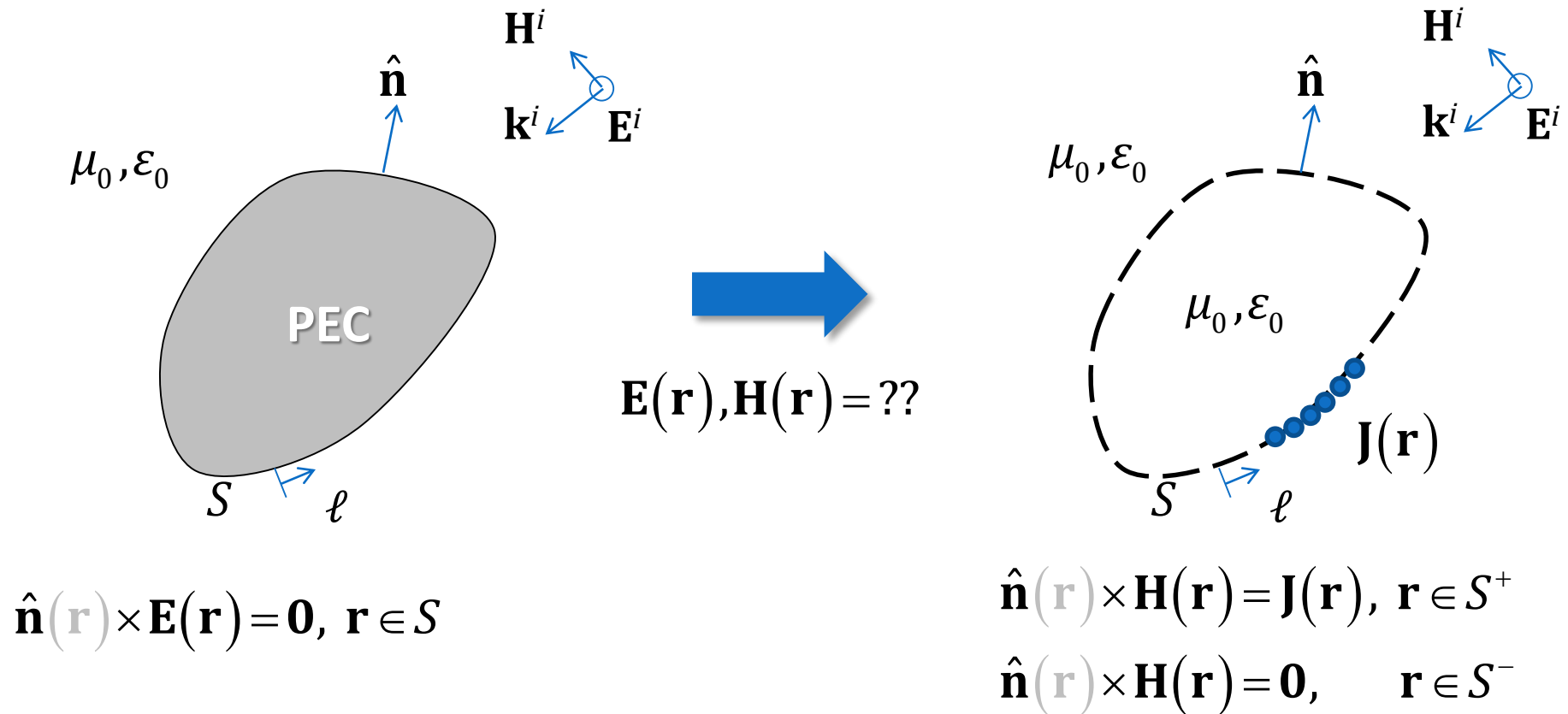
- Problem Formulation
- MoM Procedure
- Sample Results

## 2 TE Polarization

- Problem Formulation
- MoM Procedure
- Sample Results

## Problem Formulation

## Equivalent Problem



It is required to find the scattered electric and magnetic electric field everywhere due to a conducting cylinder illuminated by a plane wave.

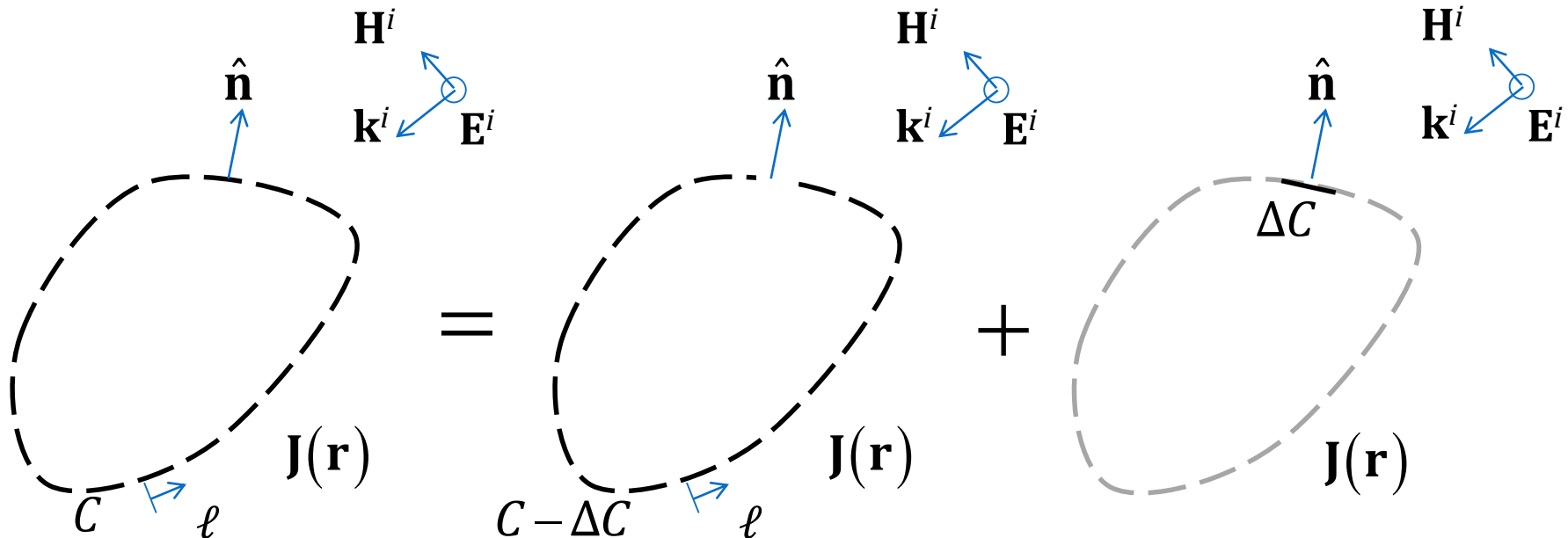
# Problem Formulation

## The Magnetic Field and Electric Current

$$\hat{\mathbf{n}} \times \mathbf{H}^+ = \hat{\mathbf{n}} \times [\mathbf{H}^{i+} + \mathbf{H}^{s+}] = \hat{\mathbf{n}} \times [\mathbf{H}^{i+} + \mathbf{H}_{ns}^{s+} + \mathbf{H}_s^{s+}] = \mathbf{J}$$

$$\hat{\mathbf{n}} \times \mathbf{H}^- = \hat{\mathbf{n}} \times [\mathbf{H}^{i-} + \mathbf{H}^{s-}] = \hat{\mathbf{n}} \times [\mathbf{H}^{i-} + \mathbf{H}_{ns}^{s-} + \mathbf{H}_s^{s-}] = \mathbf{0}$$

$$\hat{\mathbf{n}} \times [\mathbf{H}_s^{s+} - \mathbf{H}_s^{s-}] = \mathbf{J} \quad \Rightarrow \quad \boxed{\hat{\mathbf{n}} \times \mathbf{H}_s^{s\pm} = \pm \frac{1}{2} \mathbf{J}}$$



## Problem Formulation

## Maue's Magnetic Field Integral Equation

$$\hat{\mathbf{n}} \times [\mathbf{H}^i(\mathbf{r}) + \mathbf{H}_{ns}^s(\mathbf{r}; \mathbf{J})] = \frac{1}{2} \mathbf{J}(\mathbf{r}), \mathbf{r} \in S$$

$$-\frac{1}{2} \mathbf{J}(\mathbf{r}) + \frac{1}{\mu_0} \hat{\mathbf{n}} \times \nabla \times \int_{S-\Delta S} \mathbf{J}(\mathbf{r}') g(\mathbf{r}, \mathbf{r}') ds' = -\hat{\mathbf{n}} \times \mathbf{H}^i(\mathbf{r}), \mathbf{r} \in S$$

$$\nabla \times (\mathbf{J}g) = \cancel{g \nabla \times \mathbf{J}} - \mathbf{J} \times \nabla g$$

$$\nabla \times (\mathbf{J}g) = -\mathbf{J} \times \nabla g = \mathbf{J} \times \nabla' g$$

$$-\frac{1}{2} \mathbf{J}(\mathbf{r}) + \frac{1}{\mu_0} \hat{\mathbf{n}} \times \int_{S-\Delta S} \mathbf{J}(\mathbf{r}') \times \nabla' g(\mathbf{r}, \mathbf{r}') ds' = -\hat{\mathbf{n}} \times \mathbf{H}^i(\mathbf{r}), \mathbf{r} \in S$$

## Problem Formulation

## Operator Equation

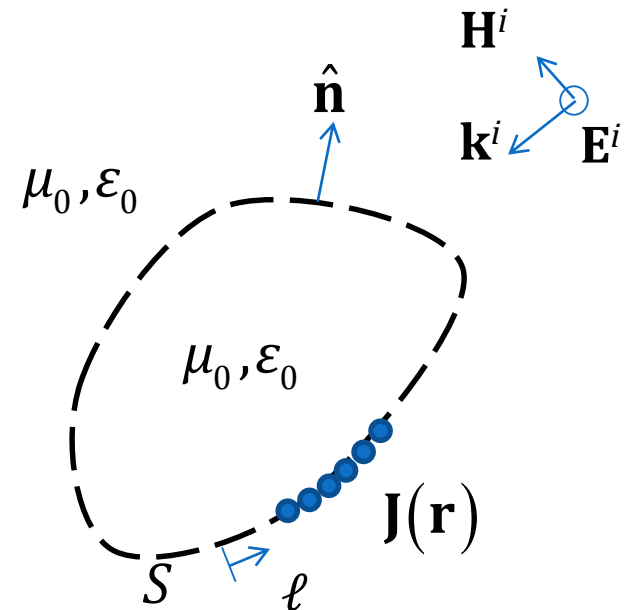
$$\mathbf{J}(\boldsymbol{\rho}) = J_z(\boldsymbol{\rho}) \hat{\mathbf{z}}$$

$$\begin{aligned} -\frac{1}{2} J_z(\boldsymbol{\rho}) + \frac{k_0}{4j} \int_{C-\Delta C} J_z(\boldsymbol{\rho}') \hat{\mathbf{n}} \cdot \hat{\mathbf{R}} H_1^{(2)}(k_0 |\boldsymbol{\rho} - \boldsymbol{\rho}'|) d\mathbf{l}' \\ = -H_1^i(\boldsymbol{\rho}), \quad \boldsymbol{\rho} \in C \end{aligned}$$

$$L\{J_z(\boldsymbol{\rho})\} = -H_1^i(\boldsymbol{\rho}), \quad \boldsymbol{\rho} \in C$$

$$H_1^i = \hat{\mathbf{l}} \cdot \mathbf{H}^i = \frac{1}{\eta_0} \hat{\mathbf{l}} \cdot \hat{\mathbf{k}}^i \times \mathbf{E}^i, \quad \boldsymbol{\rho} \in C$$

$$E_z^i(\boldsymbol{\rho}) = E_0 e^{-j\mathbf{k}^i \cdot \boldsymbol{\rho}} = E_0 e^{+jk_0(\cos\varphi^i x + \sin\varphi^i y)}$$



Based on the previous derivations, the MFIE formulation can be used only for closed objects.

## The MoM Procedure

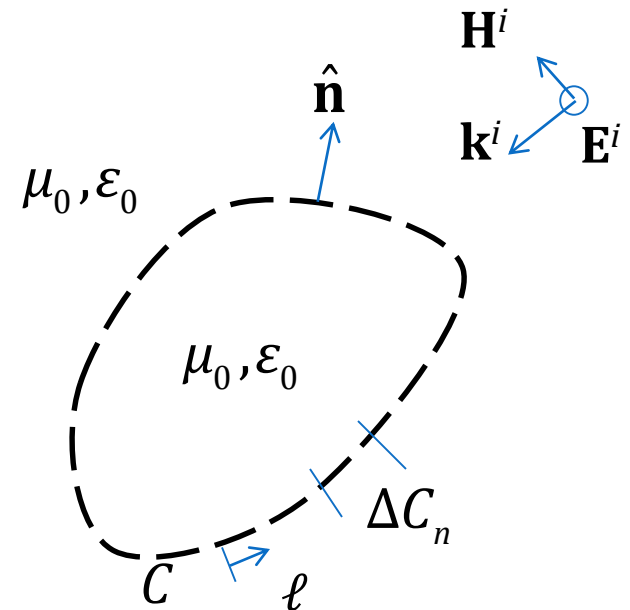
## Expansion, Testing and MoM Matrix Equation

$$\Pi_n(\boldsymbol{\rho}) = \begin{cases} 1, & \boldsymbol{\rho} \in \Delta C_n \\ 0, & \text{otherwise} \end{cases} \quad t_m(\boldsymbol{\rho}) = \delta(\boldsymbol{\rho} - \boldsymbol{\rho}_m)$$

$$\sum_{n=1}^N I_n Z_{mn} = V_m, \quad m = 1, 2, L, N$$

$$Z_{mn} = -\frac{1}{2}\delta_{mn} + \frac{k_0}{4j} \int_{\Delta C_n} \hat{\mathbf{n}}_m \cdot \hat{\mathbf{R}} H_1^{(2)}(k_0 |\boldsymbol{\rho}_m - \boldsymbol{\rho}'|) dl'$$

$$V_m = -\frac{1}{\eta_0} E_0 e^{+jk_0(\cos\varphi^i x_m + \sin\varphi^i y_m)} (\cos\varphi^i \hat{\mathbf{y}} - \sin\varphi^i \hat{\mathbf{x}}) \cdot \hat{\mathbf{l}}_m$$



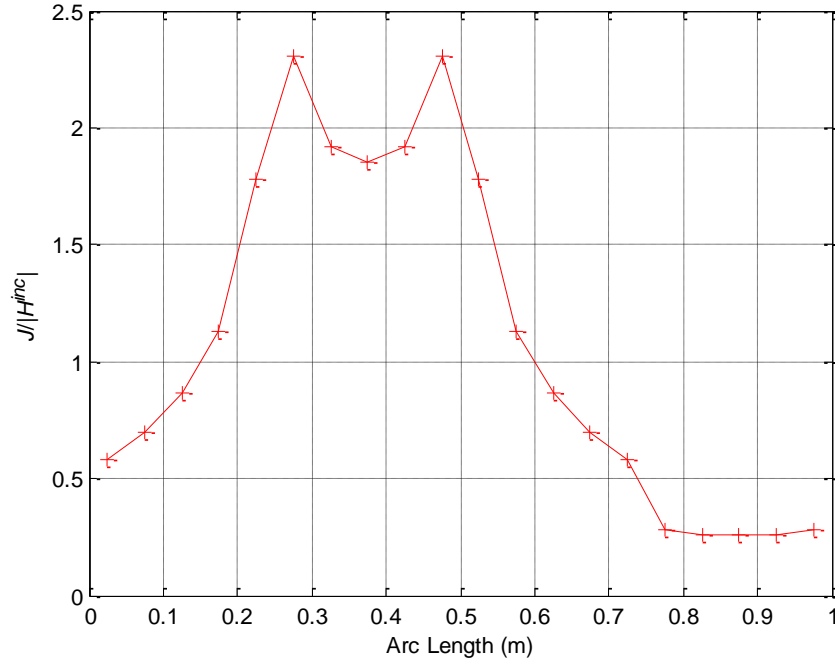
Notice that there are no singularity issues in this case. However, near-self terms may still need singularity extraction.



# Sample Results

## Sample Case 1: Quarter-Wavelength PEC Square

Normalized Current Density ( $\phi^{inc} = 0$ )



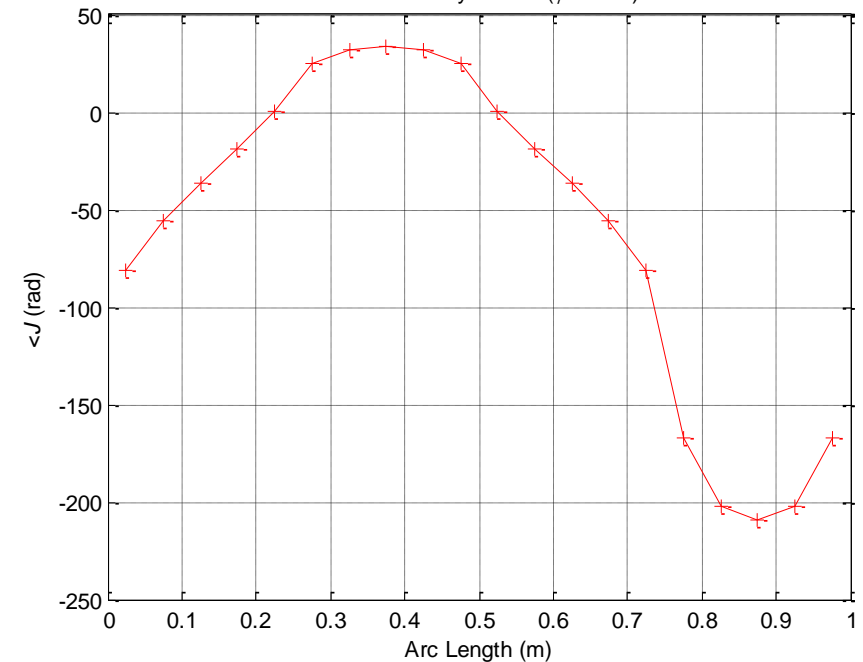
$f = 300$  MHz

$\lambda = 1$  m



← TM Plane Wave

Current Density Phase ( $\phi^{inc} = 0$ )

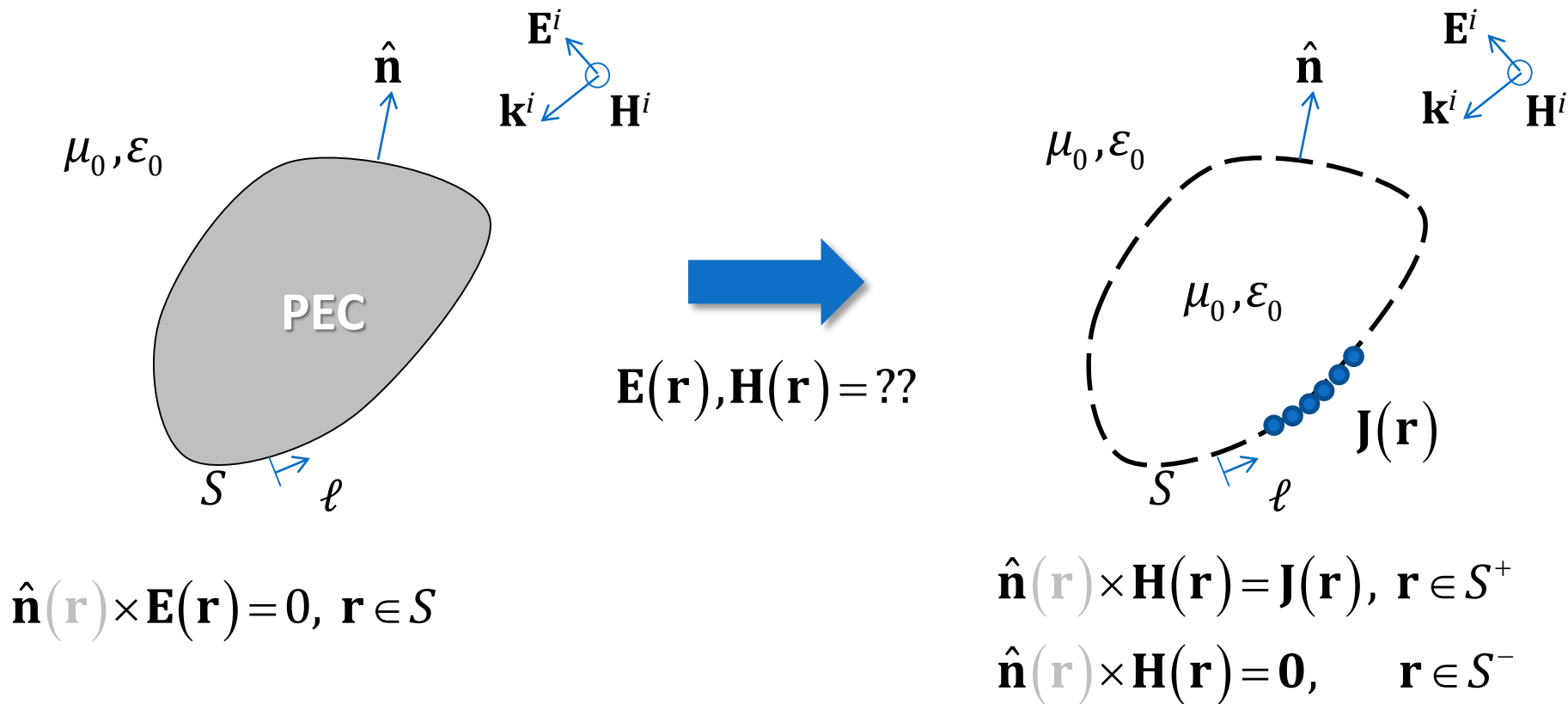


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## Problem Formulation

## Equivalent Problem



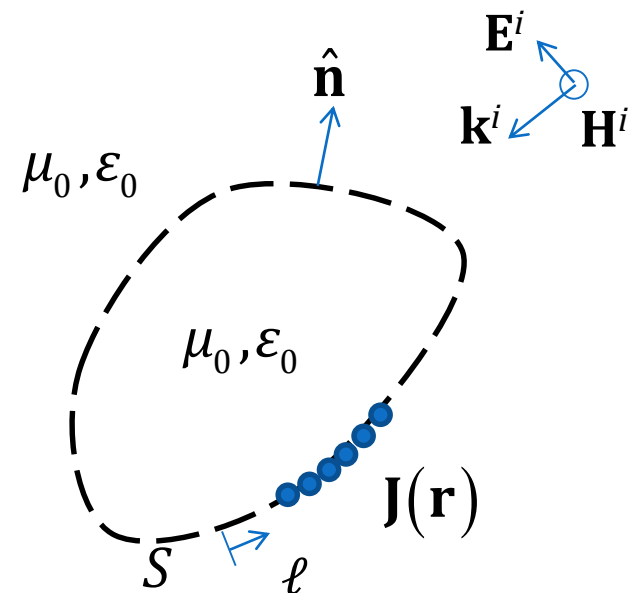
It is required to find the scattered electric and magnetic electric field everywhere due to a conducting cylinder illuminated by a plane wave.

## Problem Formulation

## Operator Equation

$$\mathbf{J}(\boldsymbol{\rho}) = J_1(\boldsymbol{\rho}) \hat{\mathbf{l}}$$

$$-\frac{1}{2} J_1(\boldsymbol{\rho}) + \frac{k_0}{4j} \int_{C-\Delta C} J_1(\boldsymbol{\rho}') \hat{\mathbf{n}}' \cdot \hat{\mathbf{R}} H_1^{(2)}(k_0 |\boldsymbol{\rho} - \boldsymbol{\rho}'|) dl' = H_z^i(\boldsymbol{\rho}), \quad \boldsymbol{\rho} \in C$$



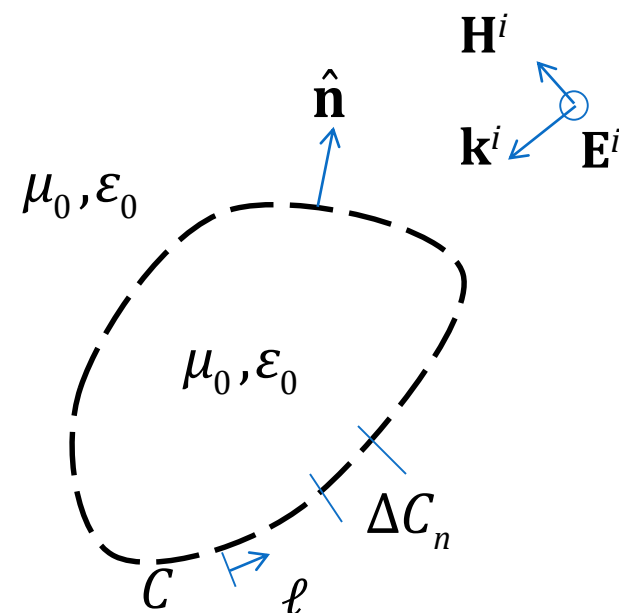
For TE polarization, the current has no axial component. Notice that a TE polarized wave can be also generated by an magnetic line source.

## The MoM Procedure

## Expansion using Vector Basis Functions

$$\mathbf{\Lambda}_n(\boldsymbol{\rho}) = \begin{cases} \hat{\mathbf{l}}_n \left( \frac{1}{\Delta l_n} + 1 \right), & \boldsymbol{\rho} \in \Delta C_n \\ \hat{\mathbf{l}}_{n+1} \left( 1 - \frac{1}{\Delta l_{n+1}} \right), & \boldsymbol{\rho} \in \Delta C_{n+1} \\ 0, & \text{otherwise} \end{cases}$$

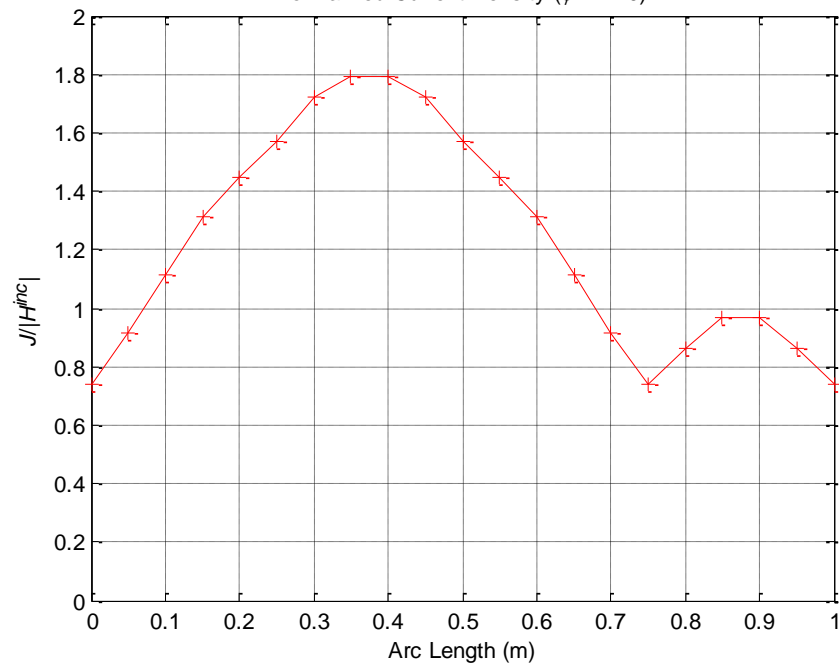
$$\mathbf{t}_n(\boldsymbol{\rho}) = \begin{cases} \hat{\mathbf{l}}_n, & \boldsymbol{\rho} \in \Delta C_n^{(2)} \\ \hat{\mathbf{l}}_{n+1}, & \boldsymbol{\rho} \in \Delta C_{n+1}^{(1)} \\ 0, & \text{otherwise} \end{cases}$$



For basis functions that span multiple segments, there should be special treatment for closed objects to account for the last unknown. (Explain)

## Sample Results

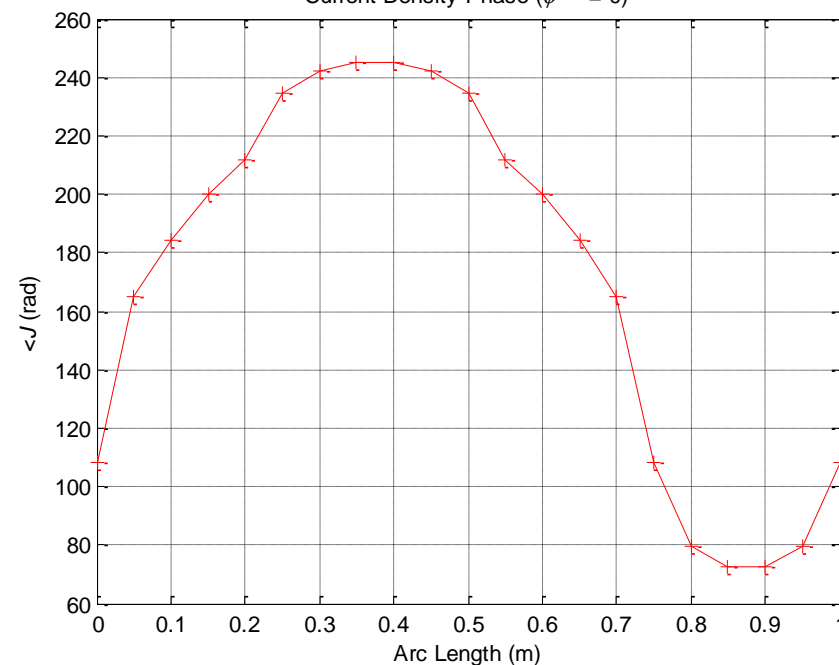
## Sample Case 1: Quarter-Wavelength PEC Square

Normalized Current Density ( $\phi^{inc} = 0$ )

$f = 300$  MHz  
 $\lambda = 1$  m



← TE Plane  
Wave

Current Density Phase ( $\phi^{inc} = 0$ )

## Assignment

## Assignment #5

Repeat Assignment #3 using MFIE formulation.

## Assignment

## Assignment #6

Repeat Assignment #5 for TE polarization.



## Conclusion

- MFIE formulation is obtained by enforcing the boundary condition on the discontinuous field component.
- The MFIE does not have singular kernels and thus needs no special treatment for self-terms.
- Comparing the EFIE and the MFIE shows that the latter yields a “better” matrix equation. However, the EFIE can handle both open and closed objects.