

# Introduction to Computational Electromagnetics

## Lecture 1: Introduction

ELC 657 – Spring 2017

Department of Electronics and Communications  
Engineering  
Faculty of Engineering – Cairo University

# Outline

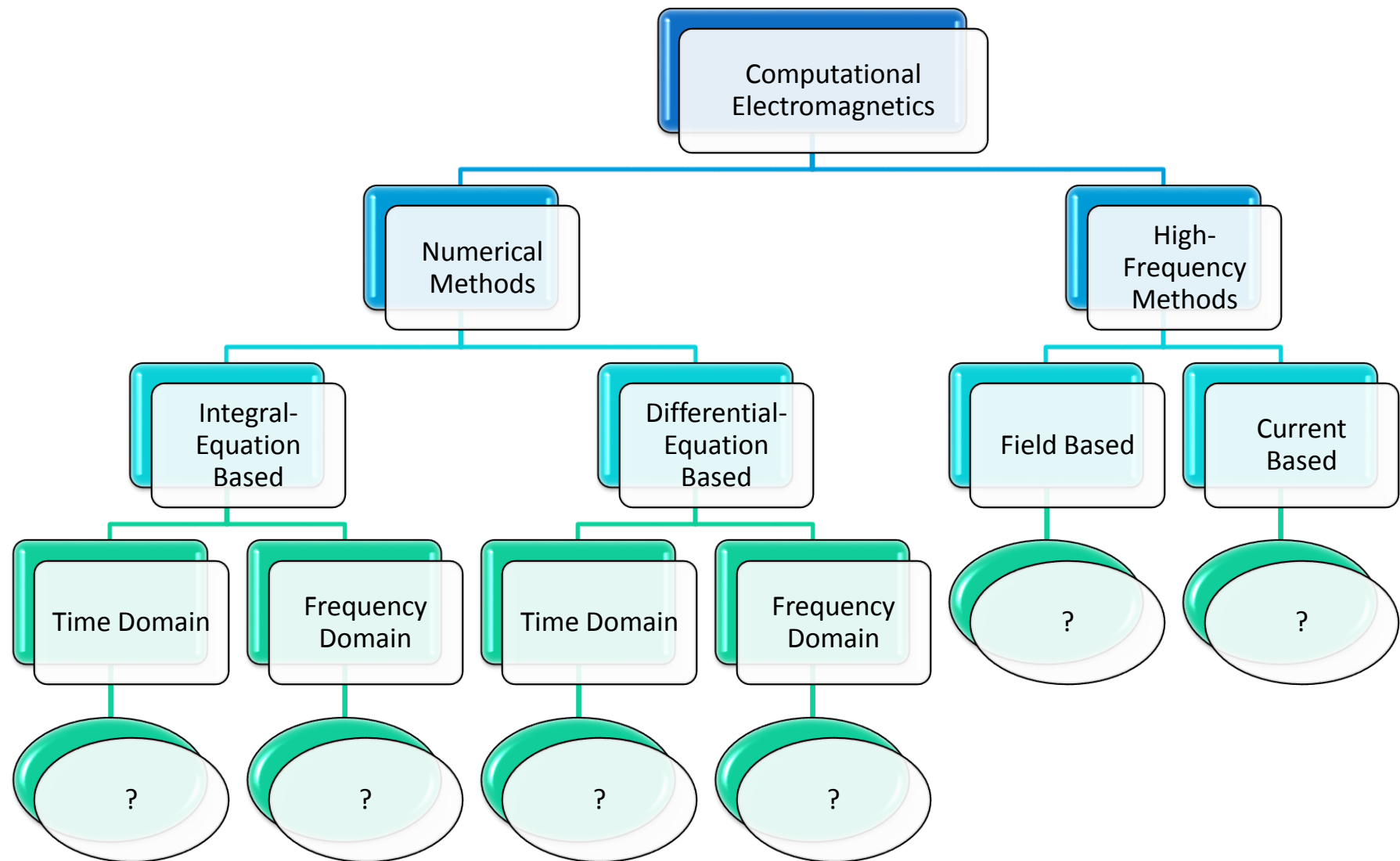
- 1 CEM Techniques
  - Classification
  - Comparison and Limitations
  - Challenges and Recent Advances
- 2 The Surface Equivalence Principle
  - Historical Background
  - Surface Integral Equation (SIE) Formulation
- 3 The Green's Function Method
  - Typical Solution Procedure
  - Known Green's Functions

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## Classification

## Classification Tree of CEM Techniques

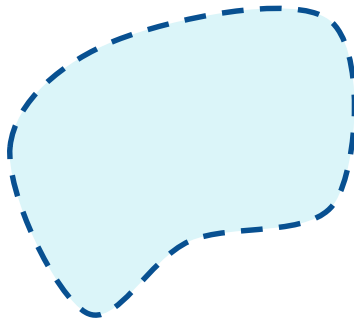


## Comparison and Limitations

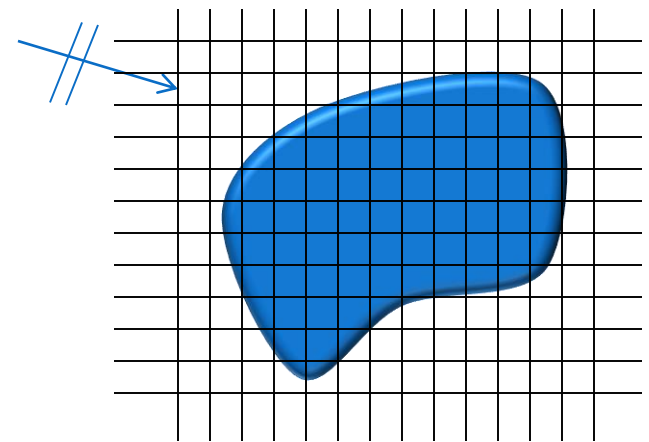
## Domain Discretization and Truncation and Memory Requirements



Partial Domain Discretization



Complete Domain Discretization



## Comparison and Limitations

## Time vs. Frequency Domain



The diagram consists of two side-by-side vertical rectangular boxes with rounded corners. Each box has a blue header bar at the top containing white text. The left box is labeled 'Time Domain' and the right box is labeled 'Frequency Domain'. The main body of each box is empty, intended for notes or diagrams comparing the two domains.

**Time Domain**

**Frequency Domain**

## Challenges and Recent Advances

## Some Challenges in CEM

- Large objects
- Distant objects
- Adaptive meshing/ Sub-meshing
- Modeling new/artificial media
- Using GPU cards
- Optimization techniques

## This Course

## Course Topics

- Introduction to CEM
  - Using the MoM! (2D PEC Electrostatic Problem)
  - Background Theory (Projection Theorem)
  - The MoM Procedure
  - The MoM applied to 2D Electromagnetic Scattering Problem  
(EFIE – MFIE – CFIE)
- 
- Introduction to Finite Differences
  - Using the FD Technique!
  - Applications to Electrostatic Problems, Transmission Lines... etc.



## This Course

## Grade Distribution

- Assignments/Presentations/...etc: 40%
- Projects: 20%
  - Project 1: General 2D PEC Electrostatic Problem
  - Project 2: General 2D PEC Electromagnetic Problem
  - Project 3: FD Solver
- Final: 40%

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*\* The projects should be written in Matlab, C, or Fortran.*

## This Course

## References

- R. F. Harrington, *Field Computation by Moment Methods*. IEEE Press, 1993.
- A. W. Glisson and Atef Elsherbeni, Lecture Notes of ELE 528 and 628. The University of Mississippi, 2002-2003.
- C. A. Balanis, *Advanced Engineering Electromagnetics*. John Wiley and Sons, 1989.
- Selected papers.

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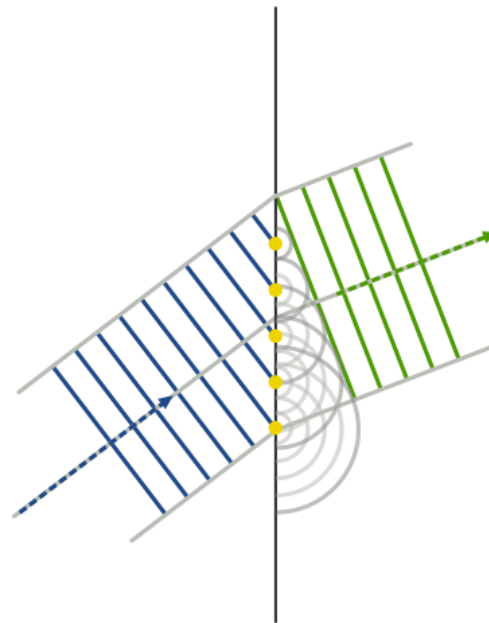
## Historical Background

# Huygens's Principle

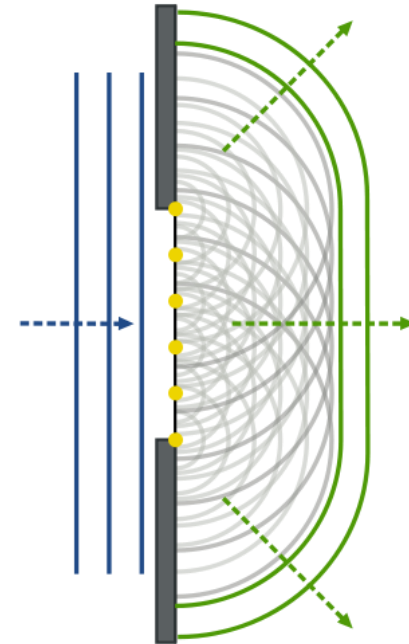


Christiaan Huygens  
(1629-1695)

Each point on the wavefront of an advancing wave (produced by some primary source) is the source of a new, secondary spherical wave.



Refraction in the manner  
of Huygens'



Diffraction in the manner  
of Huygens'

## Historical Background

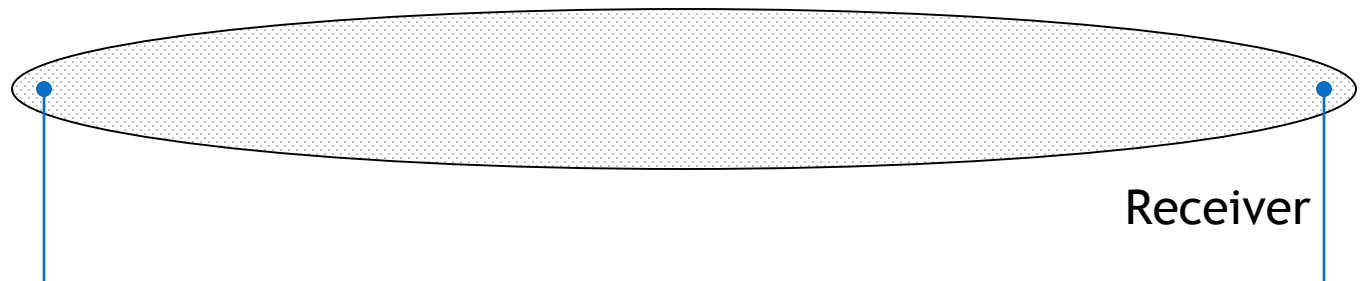
# Fresnel Zones



Augustin-Jean Fresnel  
(1788-1827)

Only a specific part of the wavefront has a significant contribution to the advancing wave, viz. the first Fresnel zone.

Transmitter



Receiver

## Historical Background

## Balanced Maxwell's Equations



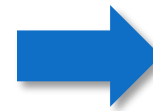
James Clerk Maxwell  
(1831-1879)

$$\nabla \times \mathbf{E} = -j\omega\mu\mathbf{H}$$

$$\nabla \times \mathbf{H} = +j\omega\epsilon\mathbf{E} + \mathbf{J}$$

$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$



$$\nabla \times \mathbf{E} = -j\omega\mu\mathbf{H} - \mathbf{M}$$

$$\nabla \times \mathbf{H} = +j\omega\epsilon\mathbf{E} + \mathbf{J}$$

$$\nabla \cdot \mathbf{D} = \rho_e$$

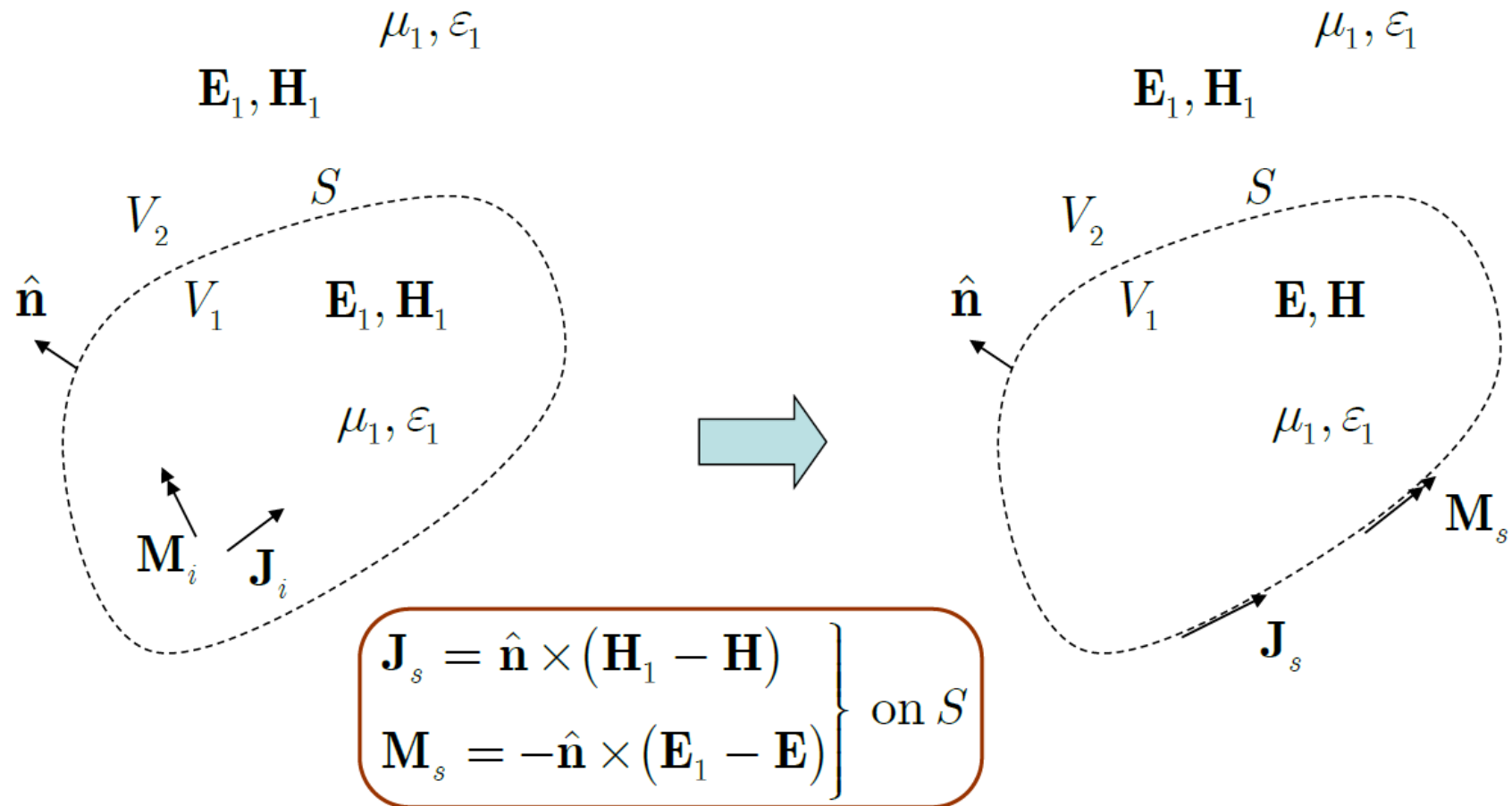
$$\nabla \cdot \mathbf{B} = \rho_m$$

$$\mathbf{E} = \underbrace{-j\omega\mathbf{A} - \nabla\Phi}_{\mathbf{E}^J} - \underbrace{\frac{1}{\epsilon}\nabla \times \mathbf{F}}_{\mathbf{E}^M}$$

$$\mathbf{H} = \underbrace{-j\omega\mathbf{F} - \nabla\Psi}_{\mathbf{H}^M} + \underbrace{\frac{1}{\mu}\nabla \times \mathbf{A}}_{\mathbf{H}^J}$$

## Historical Background

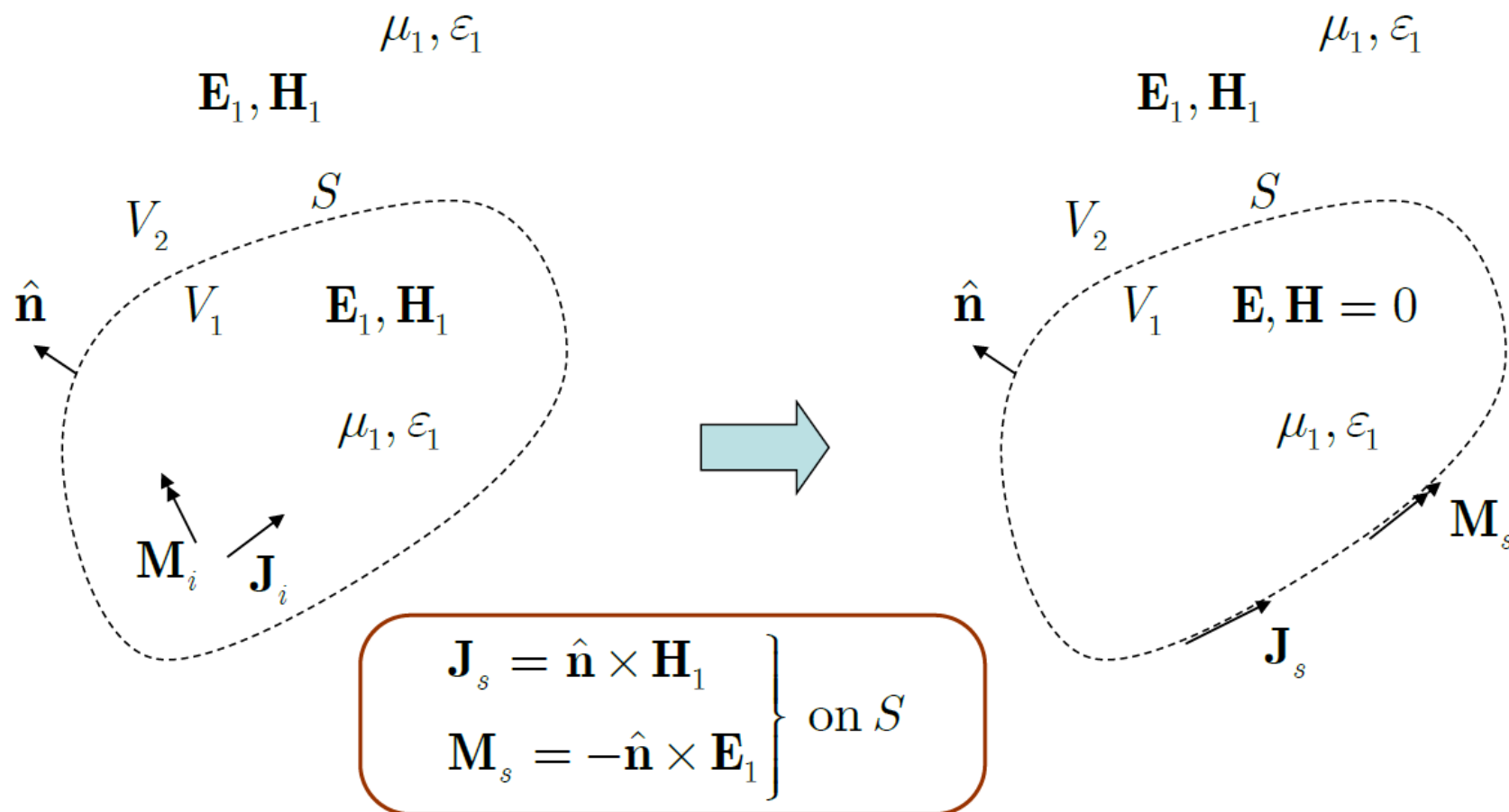
## Schelkunoff's Equivalence Principle



In Schelkunoff's equivalence, both fields need to be known on the surface, which is a redundancy according to the uniqueness theorem.

## Historical Background

## Love's Equivalence Principle

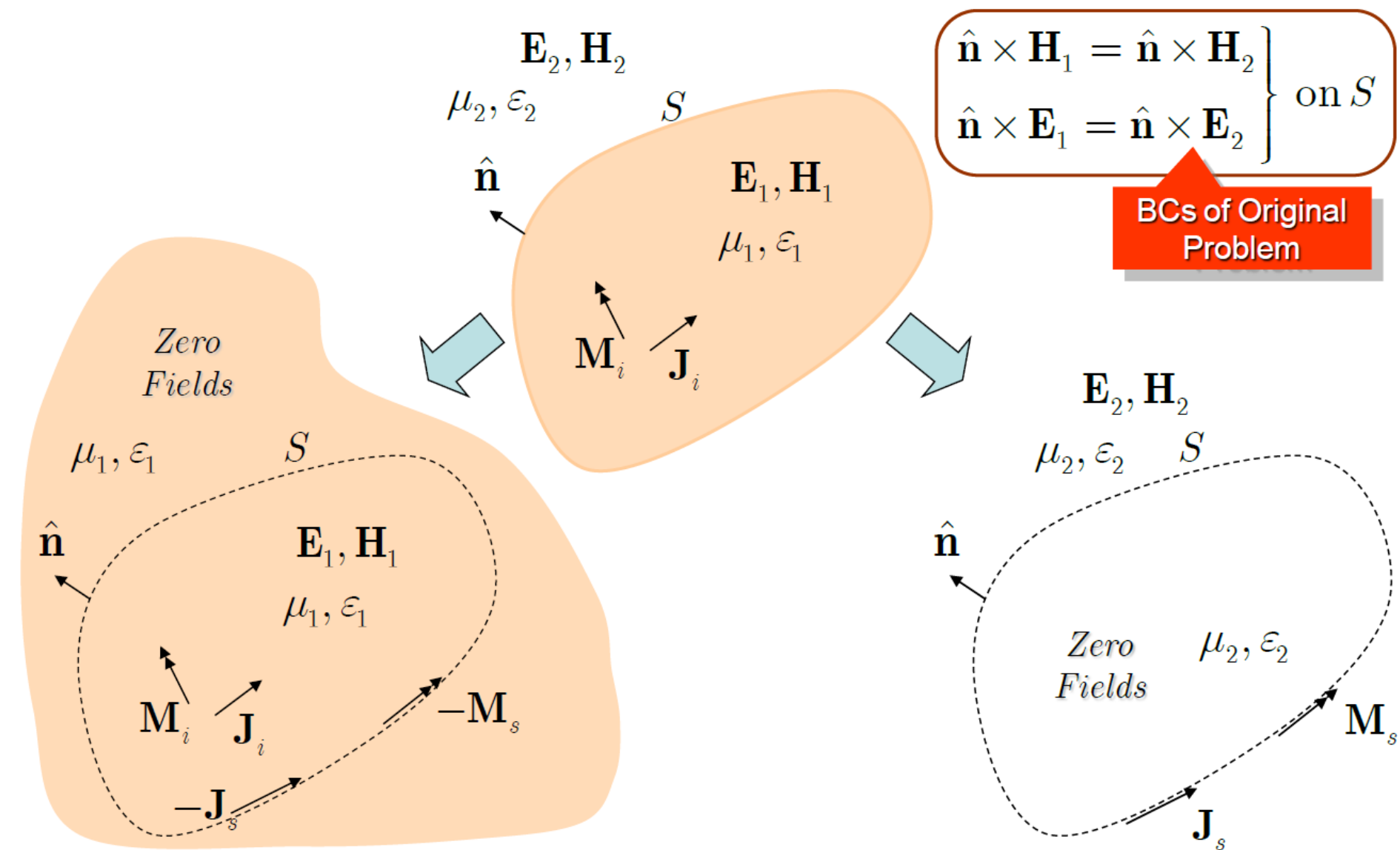


Notice that in the SE, the interior volume ( $V_1$ ) can be filled with any homogeneous medium or perfectly conducting medium. But...?



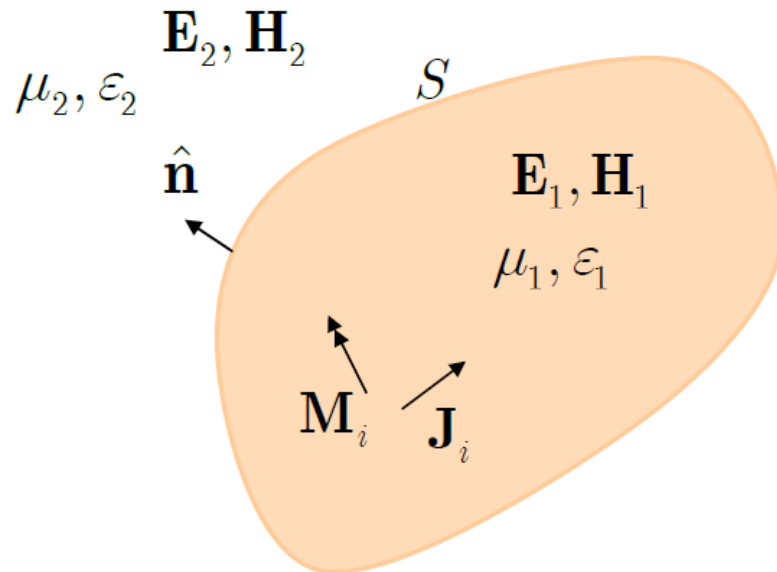
# Surface Integral Equation (SIE) Formulation

## Equivalent Problems



## Surface Integral Equation (SIE) Formulation

## Surface Integral Equations



$$\left. \begin{aligned} \hat{\mathbf{n}} \times \mathbf{H}_1 &= \hat{\mathbf{n}} \times \mathbf{H}_2 \\ \hat{\mathbf{n}} \times \mathbf{E}_1 &= \hat{\mathbf{n}} \times \mathbf{E}_2 \end{aligned} \right\} \text{ on } S$$



Boundary Conditions in  
"Operator" Form

$$\left. \begin{aligned} \hat{\mathbf{n}} \times \left[ \mathbf{E}_1^J(\mathbf{J}_i) + \mathbf{E}_1^M(\mathbf{M}_i) + \mathbf{E}_1^J(-\mathbf{J}_s) + \mathbf{E}_1^M(-\mathbf{M}_s) \right] \\ = \hat{\mathbf{n}} \times \left[ \mathbf{E}_2^J(\mathbf{J}_s) + \mathbf{E}_2^M(\mathbf{M}_s) \right] \\ \hat{\mathbf{n}} \times \left[ \mathbf{H}_1^J(\mathbf{J}_i) + \mathbf{H}_1^M(\mathbf{M}_i) + \mathbf{H}_1^J(-\mathbf{J}_s) + \mathbf{H}_1^M(-\mathbf{M}_s) \right] \\ = \hat{\mathbf{n}} \times \left[ \mathbf{H}_2^J(\mathbf{J}_s) + \mathbf{H}_2^M(\mathbf{M}_s) \right] \end{aligned} \right\} \text{ on } S$$

## Surface Integral Equation (SIE) Formulation

## The Fields as Operators

Fields due to  $\mathbf{J}$ 

$$\mathbf{E}_\nu^J(\mathbf{J}) = -j\omega\mathbf{A}_\nu(\mathbf{J}) - \nabla\Phi_\nu(\mathbf{J})$$

$$\mathbf{H}_\nu^J(\mathbf{J}) = \frac{1}{\mu_\nu} \nabla \times \mathbf{A}(\mathbf{J})$$

$$\mathbf{A}_\nu(\mathbf{J}) = \mu_\nu \iiint_V g_\nu(\mathbf{r}, \mathbf{r}') \mathbf{J}(\mathbf{r}') dv'$$

$$\Phi_\nu(\mathbf{J}) = \frac{j}{\omega\epsilon_\nu} \iiint_V g_\nu(\mathbf{r}, \mathbf{r}') \nabla' \cdot \mathbf{J}(\mathbf{r}') dv'$$

Fields due to  $\mathbf{M}$ 

$$\mathbf{H}_\nu^M(\mathbf{M}) = -j\omega\mathbf{F}_\nu(\mathbf{M}) - \nabla\Psi_\nu(\mathbf{M})$$

$$\mathbf{E}_\nu^M(\mathbf{M}) = -\frac{1}{\epsilon_\nu} \nabla \times \mathbf{F}_\nu(\mathbf{M})$$

$$\mathbf{F}_\nu(\mathbf{M}) = \epsilon_\nu \iiint_V g_\nu(\mathbf{r}, \mathbf{r}') \mathbf{M}(\mathbf{r}') dv'$$

$$\Psi_\nu(\mathbf{M}) = \frac{j}{\omega\mu_\nu} \iiint_V g_\nu(\mathbf{r}, \mathbf{r}') \nabla' \cdot \mathbf{M}(\mathbf{r}') dv'$$

Free-Space Green's  
Function

$$g_\nu(\mathbf{r}, \mathbf{r}') = \frac{e^{-jk_\nu|\mathbf{r}-\mathbf{r}'|}}{4\pi|\mathbf{r}-\mathbf{r}'|}$$

$$k_\nu = \omega\sqrt{\mu_\nu\epsilon_\nu}$$

## Surface Integral Equation (SIE) Formulation

## The SIE Formulation and the MoM

The MoM is widely used in ***Surface Integral Equation*** formulations, where unknowns are placed only over the surfaces/boundaries of homogeneous regions.

- ➡ No need to discretize all space
- ➡ Reduced number of unknowns required

The MoM is easily applied to objects in unbounded regions.

- ➡ The radiation condition is automatically satisfied
- ➡ No “absorbing boundary” is required

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## Typical Solution Procedure

## Solution in Terms of Green's Function: 3D EM Problem

$$\begin{array}{cc}
 (\nabla^2 + k^2)\psi(\mathbf{r}) = s(\mathbf{r}) & (\nabla^2 + k^2)g(\mathbf{r}, \mathbf{r}') = -\delta(\mathbf{r} - \mathbf{r}') \\
 \times g & \times (-\psi) \\
 & +
 \end{array}$$

$$g(\mathbf{r}, \mathbf{r}') \nabla^2 \psi(\mathbf{r}) - \psi(\mathbf{r}) \nabla^2 g(\mathbf{r}, \mathbf{r}') = g(\mathbf{r}, \mathbf{r}') s(\mathbf{r}) + \psi(\mathbf{r}) \delta(\mathbf{r} - \mathbf{r}')$$

$$\int_V \{ g(\mathbf{r}, \mathbf{r}') \nabla^2 \psi(\mathbf{r}) - \psi(\mathbf{r}) \nabla^2 g(\mathbf{r}, \mathbf{r}') \} dv = \int_V g(\mathbf{r}, \mathbf{r}') s(\mathbf{r}) dv + \psi(\mathbf{r}')$$

$$\int_S \left\{ g(\mathbf{r}, \mathbf{r}') \frac{\partial \psi(\mathbf{r})}{\partial n} - \psi(\mathbf{r}) \frac{\partial g(\mathbf{r}, \mathbf{r}')}{\partial n} \right\} ds = \int_V g(\mathbf{r}, \mathbf{r}') s(\mathbf{r}) dv + \psi(\mathbf{r}')$$

**Green's  
Identity**

$$\psi(\mathbf{r}') = - \int_V g(\mathbf{r}, \mathbf{r}') s(\mathbf{r}) dv$$

$$\psi(\mathbf{r}) = - \int_V g(\mathbf{r}, \mathbf{r}') s(\mathbf{r}') dv$$

## Typical Solution Procedure

## Determining the Green's Function: 3D EM Problem

$$(\nabla^2 + k^2)g(\mathbf{r}, \mathbf{r}') = -\delta(\mathbf{r} - \mathbf{r}')$$

Source at  
Origin

$$\int_V (\nabla^2 + k^2)g(r) = -\delta(r)$$

$$\int_V \nabla^2 g(r) dv + k^2 \int_V g(r) dv = -1$$

$$\int_S \nabla g(r) \cdot d\mathbf{s} + k^2 \int_V g(r) dv = -1$$

$$\lim_{a \rightarrow 0} \int_S \nabla g(r) \cdot d\mathbf{s} + k^2 \lim_{a \rightarrow 0} \int_V g(r) dv = -1$$

$$\lim_{a \rightarrow 0} \int_S A \frac{-e^{-jka} (jka + 1)}{a^2} a^2 \sin \theta d\theta d\varphi = -1$$

$$A = \frac{1}{4\pi}$$

$$(\nabla^2 + k^2)g(r) = 0, r \neq 0$$

$$g(r) = A \frac{e^{-jkr}}{r}$$

$$g(\mathbf{r}, \mathbf{r}') = \frac{e^{-jk|\mathbf{r} - \mathbf{r}'|}}{4\pi|\mathbf{r} - \mathbf{r}'|}$$

## Known Green's Functions

## Examples of Known Green's Functions

$$g_{3D}(\mathbf{r}, \mathbf{r}') = \frac{e^{-jk|\mathbf{r}-\mathbf{r}'|}}{4\pi|\mathbf{r}-\mathbf{r}'|}$$

$$g_{2D}(\boldsymbol{\rho}, \boldsymbol{\rho}') = \frac{1}{4j} H_0^{(2)}(k|\boldsymbol{\rho}-\boldsymbol{\rho}'|)$$

$$g_{1D}(z, z') = \frac{e^{-jk|z-z'|}}{2jk}$$

Other known Green's functions include half-space, waveguide, parallel plate, and layered media Green's functions.



- CEM includes several techniques, each is fit for certain problems and has its advantages and disadvantages.
- The surface equivalence principle offers a solution to EM scattering problems in terms of a set of equivalent problems and results in surface integral equations (that can be solved using the MoM).
- The Green's function is the “spatial” impulse response of the system.